ON THE INVERSE RESONANCE PROBLEM

Rudi Weikard

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50 Congreso Nacional de la Sociedad Matemática Mexicana

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Rudi Weikard (UAB)

On the inverse resonance problem

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I am reporting on joint work with

- Christer Bennewitz (Lund)
- Matthew Bledsoe (Birmingham, AL)
- Malcolm Brown (Cardiff)
- Ian Knowles (UAB)
- Marco Marletta (Cardiff)
- Sergey Naboko (St. Petersburg)
- Roman Shterenberg (UAB)

Introduction and History

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On the inverse resonance problem

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"Über eine Frage der Eigenwerttheorie" (1928):

If indeed the spectrum [of eigenvalues] defined the differential equation completely, it would be possible, for instance, to determine practically the structure of a system of atoms from the spectrum, i.e., to solve the problem which is, so to speak, reciprocal to Schrödinger's problem.

• Consider the Schrödinger equation with real potential q

$$-y'' + qy = \lambda y.$$

q is a locally integrable function on [0, b) where $0 < b \le \infty$.

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- Unless mentioned otherwise, we assume below a Dirichlet condition at zero.

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• Ambarzumian treats only a (very) special case: the Schrödinger equation on a finite interval with continuous *q*.

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- Borg (1946) showed that, in general, two sets of eigenvalues are needed to identify a potential on an interval uniquely.
- Levinson (1949) and Marchenko (1950) used different sets of data: in addition to one set of (Dirichlet) eigenvalues one needs also either Neumann data of the eigenfunctions or the norming constants of the eigenfunctions.

• Allow $b = \infty$ and q locally integrable on [0, b).

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$$m(\lambda) = A\lambda + B + \int_{\mathbb{R}} \left(\frac{1}{t-\lambda} - \frac{t}{1+t^2}\right) d\rho(t)$$

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- Gelfand-Levitan (1951): the spectral function ρ determines q uniquely.

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- Jost function: $\psi(\cdot, 0)$
- Marchenko (1955): eigenvalues, norming constants, and scattering phase $(2i\delta(k) = \overline{\psi(k,0)}/\psi(k,0))$ determine q uniquely.

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On the inverse resonance problem

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- Recall : if Im(k) > 0 and $\psi(k, 0) = 0$ then k^2 is an eigenvalue with eigenfunction $\psi(k, \cdot)$.
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- If Im(k) ≤ 0 and ψ(k, 0) = 0 then k² is a resonance. In this case ψ(k, ·) satisfies the differential equation and the boundary condition at 0 but is not square integrable.
- Both are relevant/visible in spectroscopy.

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The inverse resonance theorem

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- a and b are determined from asymptotics as k tends to ∞ along the positive imaginary axis (ψ(k,0) ~ 1 independently of q).

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On the inverse resonance problem

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- It appears this went unnoticed for more than 40 years until Korotyaev (2000) and Zworski (2001/1988) pointed it out.

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- The uniqueness theorem requires knowledge of ALL eigenvalues and resonances.
- If q is supported on [0, R], absolutely continuous on [0, R], and has a jump discontinuity at R, then the resonances are asymptotic to the curve given by

$$\operatorname{Im}(z) = -\frac{1}{R} \ln(|\operatorname{Re}(z)|) + \frac{1}{2R} \ln(|q(R)|/4).$$

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$$\operatorname{Im}(z) = -\frac{1}{R} \ln(|\operatorname{Re}(z)|) + \frac{1}{2R} \ln(|q(R)|/4).$$

• Small changes in R or q(R) produce different asymptotics.

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On the inverse resonance problem

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- Small changes in R or q(R) produce different asymptotics.
- Large resonances are physically insignificant.
- Question: How can we state (and prove) this mathematically?

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Results

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Theorem (Marletta, Shterenberg, W.; Commun. Math. Phys. 2010)

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- q, \tilde{q} are supported in [0, 1]
- $\|q\|_1, \|\tilde{q}\|_1 \le Q$

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Then

$$\sup_{x\in[0,1]}\left|\int_x^1(q-\tilde{q})dx\right|\leq f(\varepsilon,R)$$

where $f(\varepsilon, R) \to 0$ as $R \to \infty$ but $\varepsilon R^{1/6} \to 0$.

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Stability in the case of a compact interval

• Ryabushko (1983): Suppose q_0 and q are real and have zero average. Then

$$\| q - q_0 \|_{L^2} \leq C \left(\| \lambda(q) - \lambda(q_0) \|_{\ell^2} + \| \mu(q) - \mu(q_0) \|_{\ell^2}
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where *C* depends on $||q||_2$ and $||q_0||_2$.

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- McLaughlin (1988) has a similar estimate involving one spectrum and norming constants.
- Marletta and myself (2005) gave an estimate (in terms of N and ε) on

$$\left|\int_0^x (q-q_0)dt\right| \leq f(\varepsilon,N)$$

where $f(\varepsilon, N) \to 0$ as $N \to \infty$ but $\varepsilon \log N \to 0$ provided that 2N eigenvalues are known up to an error ε .

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On the inverse resonance problem

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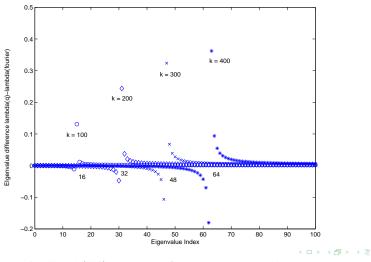
Comparison of eigenvalues

If k is large, $q(x) = \sin(kx)$, and $\tilde{q}(x) = 0$ then small eigenvalues practically coincide.

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On the inverse resonance problem

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- W., Zinchenko (Inverse Problems 2010) and Shterenberg, W., Zinchenko (Proc. Sympos. Pure Math. 2013): Uniqueness for CMV operators

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- Marletta, Naboko, Shterenberg, W. (J. Anal. Math. 2011): Stability for several classes of Jacobi operators: Spectrum is (i) all of ℝ, (ii) a half-line, or (iii) one finite interval.

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- The zeros of the reflection coefficient are also needed.

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On the inverse resonance problem

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Left-definite operators

 If q ≥ 0 but no requirement on the sign of w is made one can develop a spectral and scattering theory for

$$-y'' + qy = \lambda wy.$$

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- The inverse resonance problem for the half-line case was treated by Bledsoe, W. (J. Math. Anal. Appl., 2015)

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On the inverse resonance problem

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Outline of the proof

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The Jost solutions associated with q and \tilde{q} are related by

$$\tilde{\psi}(z,x) = \psi(z,x) + \int_{x}^{2-x} K(x,t)\psi(z,t)dt$$

where K satisfies the wave equation

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$$K_{xx}(x,t) - K_{tt}(x,t)$$
$$= (\tilde{q}(x) - q(t))K(x,t)$$

with the boundary conditions:

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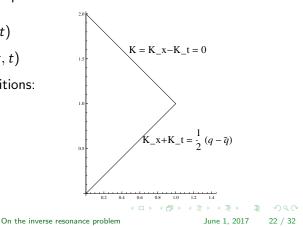
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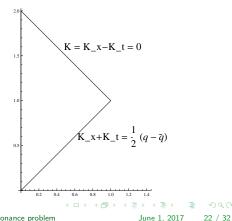
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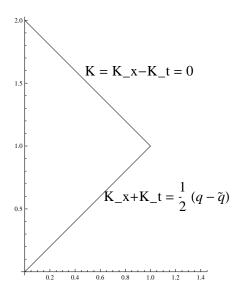
with the boundary conditions:

We need to estimate $K(x,x) = \frac{1}{2} \int_{x}^{1} (\tilde{q}(s) - q(s)) ds.$



On the inverse resonance problem

The wave equation



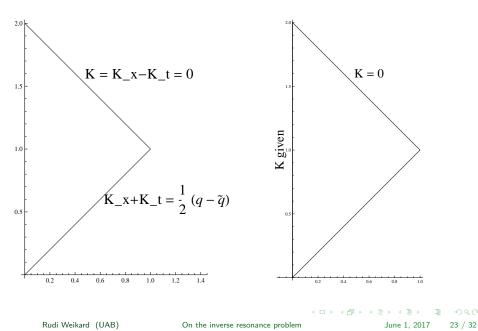
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The wave equation



Solving the wave equation

The wave equation may be solved uniquely knowing K(0, t), $0 \le t \le 2$ and the fact that K(x, 2 - x) = 0. Iteration:

$$K_0(x,t) = K(0,x+t)$$

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$$K_{n+1}(x,t) = \int_{(t+x)/2}^{1} \int_{(t-x)/2}^{(t+x)/2} (q(\alpha+\beta) - \tilde{q}(\alpha-\beta)) K_n(\alpha-\beta,\alpha+\beta) d\beta d\alpha$$

$$K(x,t) = \sum_{n=0}^{\infty} K_n(x,t)$$

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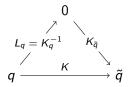
We need to estimate K(0, t).

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On the inverse resonance problem

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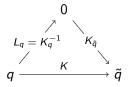
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On the inverse resonance problem

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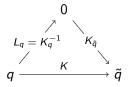
$$\mathcal{K}(0,t)=(\mathcal{K}_{\widetilde{q}}-\mathcal{K}_q)(0,t)+\int_0^t(\mathcal{K}_{\widetilde{q}}-\mathcal{K}_q)(0,s)L_q(s,t)ds$$

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$$\psi(z,0) = 1 + \int_0^2 K_q(0,t) e^{izt} dt$$

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We need to estimate $(\tilde{\psi} - \psi)(z, 0)$.

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• If f is entire of growth order at most one, then

$$f(z) = z^{n_0} \mathrm{e}^{a+bz} \prod_{n=1}^{\infty} (1-z/z_n) \mathrm{e}^{z/z_n}.$$

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• If the first $N(R) \sim 2eR$ zeros coincide

$$\frac{\psi(z,0)}{\tilde{\psi}(z,0)} = e^{(a-\tilde{a})z+b-\tilde{b}} \frac{\Pi(R,z)}{\tilde{\Pi}(R,z)}$$

where

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• $|\Pi(R,z) - 1| \le C|z|^2/R$ when 2|z| < R.

- This provides an estimate for $|z| < R^{1/3}$: $\psi(z,0)/\tilde{\psi}(z,0) \approx 1$ and hence

$$|\psi(z,0)- ilde{\psi}(z,0)|\leq CR^{-1/3}.$$

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If |z| is small: $E(z) = (1-z)e^z \approx 1-z^2$, In fact, $|\log E(z)| \le 2|z|^2$

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• $(K_{\tilde{q}} - K_q)(0, t) = h(t) + \frac{1}{2} \int_{t/2}^{1} (\tilde{q} - q)$ where h, h' is AC on [0, 2].

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- $(K_{\tilde{q}} K_q)(0, t) = h(t) + \frac{1}{2} \int_{t/2}^{1} (\tilde{q} q)$ where h, h' is AC on [0, 2].
- Integration by parts in $\tilde{\psi}(z,0) \psi(z,0) = \int_0^2 (K_{\tilde{q}} K_q)(0,t) e^{izt} dt$ gives

$$\tilde{\psi}(z,0)-\psi(z,0)=\frac{i}{z}(K_{\tilde{q}}-K_{q})(0,0)-\frac{i}{4z}\hat{G}(z)$$

where

$$\hat{G}(z) = \int_0^2 \left((\tilde{q}-q)(t/2) - 4h'(t) \right) \mathrm{e}^{izt} dt.$$

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• To deal with the last term one needs the Hausdorff-Young inequality: for 1

$$\|\hat{G}\|_q \leq rac{p^{1/(2p)}}{q^{1/(2q)}} \|G\|_p.$$

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• Here one needs the assumption that $\tilde{q} - q$ be in L^p .

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Open problem

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Two (or more) spectral bands

Suppose $q(x) = -2\wp(x + \omega)$. The spectrum of the associated Schrödinger operator has only absolutely continuous spectrum with one gap. All solutions of $-y'' + qy = \lambda y$ are explicitly known. The inverse of the map $\wp(z) = \lambda$ maps the energy (λ) plane to a parallelogram (the fundamental domain of \wp) in a one-to-two fashion. Compactly supported perturbations do not change the essential spectrum but introduce eigenvalues and resonances

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Thank you for your attention!

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