

The Sets of ground states,  $\mathcal{S}_0$  and thermal equilibrium states at inverse temperature  $\beta \in [0, \infty)$ ,  $\mathcal{S}(\beta)$ . ①

$(\mathcal{A}, \{\tau_t\}_{t \in \mathbb{R}})$  a QSS with a SCOG  
 $\uparrow$   
 on  $(\Gamma, d)$  of automorphisms  
 $\tau_t = e^{it\delta}$ ,  $\delta = \lim_{A \uparrow \Gamma} [H_A, \cdot]$

Def<sup>n</sup> (i) a state  $\omega$  on  $\mathcal{A}$  is a ground state iff  $\omega(A^* \delta(A)) \geq 0$ , for all  $A \in \mathcal{A}_{loc}$

(ii)  $\omega$  is an eq. state at  $\beta^{-1}$ , iff  
 $\omega(A^* \delta(A)) \geq \beta^{-1} \omega(A^* A) \log \left( \frac{\omega(A^* A)}{\omega(AA^*)} \right)$   
 for all  $A \in \mathcal{A}_{loc}$ .

•  $\mathcal{S}_0$  and  $\mathcal{S}_\beta$  are convex sets

• extreme points of  $\mathcal{S}_0$  are pure states  
 $(\mathcal{H}_\omega, \pi_\omega, \Omega_\omega)$  in the GNS triple of  $\omega$

$\rightarrow$  and have  $\overline{\pi_\omega(\mathcal{A})}^{\text{weak}} = \mathcal{B}(\mathcal{H}_\omega)$

• the extreme points are factor states:  
 meaning center of  $\overline{\pi_\omega(\mathcal{A})}^{\text{w}} = \mathbb{C} \mathbb{1}$ .

•  $\mathcal{A}^{(B)}$  is a simplex

(2)

local symmetries / gauge symmetries

Suppose  
 $\forall x \in \Gamma$ , you have  $U_x \in \mathcal{B}(\mathcal{H}_x)$  unitary

$$\forall \Lambda \subseteq \Gamma, \alpha^{(\Lambda)}(A) = \left( \bigotimes_{x \in \Lambda} U_x^* \right) A \left( \bigotimes_{x \in \Lambda} U_x \right)$$

for each  $A \in \mathcal{A}_\Lambda$

$\alpha = \lim_{\Lambda \uparrow \Gamma} \alpha^{(\Lambda)}$  is an automorphism of  $\mathcal{A}$ .

If  $G$  is a group and  $g \mapsto U_x(g)$  is a unitary rep. of  $G$ , then this construction yields a rep  $g \mapsto \alpha_g$  of  $G$  on  $\mathcal{A}$ .

$\alpha$  is a symmetry of  $\{\tau_t\}_{t \in \mathbb{R}}$  if  $\tau_t \circ \alpha = \alpha \circ \tau_t, \forall t \in \mathbb{R}$

Sufficient condition is that

$$[H_\Lambda, \bigotimes_{x \in \Lambda} \alpha_x] = 0$$

for each  $\Lambda$ .

# Space symmetries/Point symmetries ③

$\Gamma = \mathbb{Z}^d$ , acts on itself as translations

$\forall x$  the mapping  $y \mapsto x+y$

Assuming  $\mathcal{A}_x \cong \mathbb{C}^d$  the same, for all  $x \in \mathbb{Z}^d$

then  $\forall x \in \mathbb{Z}^d$  define  $T_x$  on  $\mathcal{A}$  by extension of

$$A \in \mathcal{A}_{\{y\}} \mapsto T_x(A) \in \mathcal{A}_{\{x+y\}}$$

$$\dots \mathbb{1} \otimes \sigma^i \otimes \mathbb{1} \rightarrow \mathbb{1} \otimes \dots \otimes \sigma^i \otimes \dots$$

$\uparrow$   $\uparrow$   
 $y$   $x+y$

An interaction  $\Phi$  is translation invariant if

$$\Phi(x+z) = T_x(\Phi(x))$$

$$\Rightarrow T_t \circ T_x = T_x \circ T_t$$

- ~~if~~  $\forall \omega \in \mathcal{L}_0$  or  $\mathcal{L}^{(B)}$  are  $\alpha$ -invariant  
i.e.  $\omega \circ \alpha = \omega$  then the symmetry is unbroken
- ~~Use it in~~  $E$  (if it is (spontaneously) broken,

(Proof on board for g.s.'s)

(4)

$$\begin{aligned}\omega \circ \alpha(A^* \delta(A)) &= \omega(\alpha(A^* \delta(A))) \\ &= \omega(\alpha(A^*) \alpha(\delta(A))) \\ &= \omega(\alpha(A^*) \delta(\alpha(A))) \\ &= \omega(\alpha(A)^* \delta(\alpha(A))) \geq 0.\end{aligned}$$

## General Theorem

①  $\Gamma = \mathbb{Z}$  and short range  $\mathbb{I}$

~~$\mathcal{S}^{(B)} = \{\omega\}$~~ ,  $\exists!$  equilibrium state

$\forall \beta \in [0, \infty)$ . Araki (1968)

② Mermin-Wagner-Hohenberg  
1966                      1967

Frohlich & Pfister 1981

Continuous symmetries are unbroken for  $\beta < \infty$ ,  
 $d \leq 2$

and not too-long-range  
interactions

③ The Goldstone-Theorem (Ladav-Perez-Wroński) <sup>②</sup>  
JSP 1981

For translation invariant models on  $\mathbb{Z}^d$ ,  
ground states: continuous symmetry breaking  
implies gapless excitation spectrum

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excitation spectrum =  $\text{spec}(H_\omega) \setminus \{0\}$

"spectral gap" above the g.s.

$$\gamma = \sup \{ \delta > 0 \mid (0, \delta) \cap \text{spec}(H_\omega) = \emptyset \}$$

( $\gamma = 0$  if  $\{ \} = \emptyset$ )

Equilibrium states continuous symmetry breaking  
 $\Rightarrow$  slow decay of correlations

④ Exponential Clustering Theorem  
(2006 Hastings & Koma)

If  $\delta_0 = \{ \omega \}$ ,  $\Phi$  has exponential decay  
and  $\gamma > 0$ , then all correlations  
in  $\omega$  decay exponentially in space.

# Frustration free spin chains

(7)

$$\Gamma = [1, \infty) \subseteq \mathbb{Z}, \quad \mathcal{H}_x \cong \mathbb{C}^d \quad \text{for all } x \in \Gamma$$

Finite range interactions  $\overline{\Phi}([x, x+R])$   
 $\in \mathcal{A}_{[x, x+R]}$

$R=1$ : nearest neighbor interactions

$$H_{[1, N]} = \sum_{x=1}^{N-R} h_x$$

call this  $h_x$

WLOG assume  $\therefore \inf \text{spec}(h_x) = 0$

Frustration free = FF

FF  $\inf \text{spec}(H_{[1, N]}) = 0$  for all  $N \geq R$

$\Leftrightarrow \ker(H_{[1, N]}) \neq \{0\} \quad \forall N \geq R+1$

$$G_x = \ker(h_x)$$

$$\Leftrightarrow \bigwedge_x (\mathbb{C}^d)^{\otimes x-1} \otimes G_x \otimes (\mathbb{C}^d)^{\otimes N-R-x-1} \neq \{0\}$$

$\uparrow$   
 $\mathcal{H}_{[x, x+R]}$

It has dimension  $> 0$ .

Let  $\omega_N$  be a ground state of  $H_{\Gamma, \omega_N}$  ⑧

$\Rightarrow$  any weak limit point of  $\{\omega_N\}$ ,

$\omega$  will be a g.s. of  $\mathcal{A}_{\Gamma, \omega}$ ,

and  $\omega(h_x) = 0$ , for each  $x \in \Gamma$ .

Cuntz algebra  $\mathcal{O}(d)$  Is the unique

$C^*$ -algebra with  $\mathbb{1}, S_1, \dots, S_d,$

$S_1^*, \dots, S_d^*$ ,

$$\sum_{\alpha=1}^d S_{\alpha}^* S_{\alpha} = \mathbb{1}, \quad S_{\alpha} S_{\beta}^* = \delta_{\alpha\beta} \mathbb{1} \text{ for } \alpha, \beta \in \{1, \dots, d\}$$

$$\sigma: \mathcal{A}_{\Gamma, \omega} \rightarrow \mathcal{O}(d)$$

$\{|1\rangle, \dots, |d\rangle\}$  o.n. basis of  $\mathbb{C}^d$

$$\sigma(|\alpha_1 \dots \alpha_n\rangle \langle \beta_1 \dots \beta_n|) = S_{\alpha_1}^* \dots S_{\alpha_n}^* S_{\beta_n} \dots S_{\beta_1}$$

$$\underline{\alpha} = (\alpha_1, \dots, \alpha_n)$$

$$S_{\underline{\alpha}} = S_{\alpha_n} \dots S_{\alpha_1}$$

$$\sigma(|\underline{\alpha}\rangle \langle \underline{\beta}|) = S_{\underline{\alpha}}^* S_{\underline{\beta}}$$

$$|\alpha\rangle\langle\beta| \cdot |\gamma\rangle\langle\delta| = \delta_{\beta\gamma} |\alpha\rangle\langle\delta|$$

$$S_\alpha^* S_\beta S_\gamma^* S_\delta = \delta_{\beta\gamma} S_\alpha^* S_\delta$$

$$\sum_\alpha |\alpha\rangle\langle\alpha| = \mathbb{1}$$

Let  $\omega$  be a pure zero-energy ground state i.e.  $\omega(Hx) = 0, \forall x \in \mathbb{C}^d$

Define a state  $\tilde{\omega}$  on  $\mathcal{O}(d)$

$$\tilde{\omega}(S_\alpha^* S_\beta) = \omega(|\alpha\rangle\langle\beta| \otimes \mathbb{1} \dots)$$

$$\tilde{\omega}(\prod_i S_i^\#) = 0 \quad \text{if } \# S^* \neq \# S \text{'s}$$

$\{|1\rangle, \dots, |d\rangle\}$  is o.n. basis of  $\mathbb{C}^d$

Thm 1 For any state  $\omega$  on  $\mathcal{A}[\mathbb{Z}, \infty)$ ,

$\exists$  a Hilbert space  $\mathcal{H}$ , a unit  $\Omega \in \mathcal{H}$

and  $v_1, \dots, v_d \in \mathcal{B}(\mathcal{H})$  s.t.

(i)  $\sum_{\alpha=1}^d v_\alpha^* v_\alpha = \mathbb{1}$ , (ii)  $\mathcal{H} = \text{span} \{ v_{\alpha_1} \dots v_{\alpha_n} \Omega \mid \alpha_i, \dots, \alpha_n \}$

(iii)  $\omega(|\alpha\rangle\langle\beta|) = \langle v_\alpha \Omega \mid v_\beta \Omega \rangle$



$$\mathcal{K} \subseteq \mathcal{H}_{\omega}, \quad \nu_{\alpha} = \pi_{\omega}(S_{\alpha})|_{\mathcal{K}}$$

$$\mathcal{K} = \overline{\{ \pi(S_{\alpha})\Omega : n \geq 0, \alpha_1, \dots, \alpha_n \}} \\ \subseteq \mathcal{H}_{\omega}$$

Thm 2 Let  $\omega$  be a pure zero-energy g.s. on  $\mathcal{B}(\mathcal{K})$

with the  $\mathcal{K}$  and  $\{ \nu_{\alpha} \}_{\alpha=1, \dots, d}$  of Thm 1,

iff (i)  $\forall \psi \in G_{\mathcal{K}}^{\perp}, \psi = \sum_{\alpha_0 \dots \alpha_k} \psi_{\alpha_0 \dots \alpha_k} |\alpha_0 \dots \alpha_k\rangle$   
↑  
coeff.'s

we have the condition

$$\sum_{\alpha_0 \dots \alpha_k} \overline{\psi_{\alpha_0 \dots \alpha_k}} \nu_{\alpha_0} \dots \nu_{\alpha_k} = 0$$

(ii) for any  $X \in \mathcal{B}(\mathcal{K}), \sum_{\alpha} \nu_{\alpha}^* X \nu_{\alpha} = X$   
implies  $X = \lambda \mathbb{1}, \lambda \in \mathbb{C}.$

Proof  $\forall \psi \in G_x^+$ ,

(1)

$$0 = \omega(\mathbb{1}_{[1, \dots, x-1]} \otimes |\psi\rangle\langle\psi| \otimes \mathbb{1}_{[x+R+1, \infty)})$$

$$0 = \sum_{\alpha, \beta} \psi_{\alpha} \bar{\psi}_{\beta} \omega(|\alpha, \alpha\rangle\langle\beta, \beta|)$$

$\alpha, \beta, \beta$   
 $\uparrow \quad \uparrow \uparrow$   
 $x-1$       $R$  indices  
indices    each

$$0 = \sum_{\gamma} \langle v_{\gamma} \Omega, \left( \sum_{\alpha} \psi_{\alpha} v_{\alpha}^* \right) \left( \sum_{\beta} \bar{\psi}_{\beta} v_{\beta} \right) v_{\gamma} \Omega \rangle$$

$$= \sum_{\gamma} \langle v_{\gamma} \Omega, \left( \sum_{\alpha} \bar{\psi}_{\alpha} v_{\alpha} \right)^* \left( \sum_{\beta} \bar{\psi}_{\beta} v_{\beta} \right) v_{\gamma} \Omega \rangle$$

$$\Rightarrow \left( \sum_{\alpha} \bar{\psi}_{\alpha} v_{\alpha} \right)^* \left( \sum_{\beta} \bar{\psi}_{\beta} v_{\beta} \right) = 0$$

$$\Rightarrow \sum_{\beta} \bar{\psi}_{\beta} v_{\beta} = 0 \quad \text{so (i) is established}$$

(Note all arrow can be reversed

for iff not just only if)

To prove (ii), Define  $\hat{E}(X) = \sum_{\alpha=1}^d v_{\alpha}^* X v_{\alpha}$  (12)

Suppose  $X \neq \lambda \mathbb{1}$  for any  $\lambda \in \mathbb{C}$  and  $\hat{E}(X) = X$

$$\Rightarrow \hat{E}(X^*) = X^* \text{ and } \hat{E}(\mathbb{1}) = \mathbb{1}$$

$$\exists Y = c_1 X + c_2 X^* + c_3 \mathbb{1} \text{ s.t. } 0 \leq y \leq \mathbb{1} \text{ and } y \neq \lambda \mathbb{1}$$

Make a new state  $\omega_Y$ :

$$\omega_Y(|\alpha\rangle\langle\beta|) = \langle v_{\alpha} \Omega, Y v_{\beta} \Omega \rangle$$

Can also make  $\omega_{\mathbb{1}-Y}$ .

Compatibility means

$$\forall \underline{\alpha} = (\alpha_1, \dots, \alpha_n), \forall \underline{\beta} = (\beta_1, \dots, \beta_n)$$

$$\sum_{\underline{\gamma}} \omega_Y(|\underline{\alpha}\gamma\rangle\langle\underline{\beta}\gamma|) = \omega(|\underline{\alpha}\rangle\langle\underline{\beta}|)$$

