

Localization for Random Block Operators Related to the XY Spin Chain

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The XY Spin Chain: The Model

- The anisotropic XY spin chain is given by the self-adjoint Hamiltonian

$$H_n = \sum_{j=1}^{n-1} \mu_j [(1 + \gamma_j) \sigma_j^x \sigma_{j+1}^x + (1 - \gamma_j) \sigma_j^y \sigma_{j+1}^y] + \sum_{j=1}^n \nu_j \sigma_j^z$$

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$$\sigma^x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma^y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma^z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

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- $M_j := I \otimes \cdots \otimes I \otimes M \otimes I \otimes \cdots \otimes I$ (nontrivial in j th component)

Random Block Operators

- Via the Jordan-Wigner transform, one reduces the problem to proving dynamical localization for the random block operator

$$\hat{M}_n = \begin{pmatrix} A_n & B_n \\ -B_n & -A_n \end{pmatrix}$$

where

$$A_n = \begin{pmatrix} \nu_1 & -\mu_1 & & & \\ -\mu_1 & \ddots & & \ddots & \\ & \ddots & \ddots & & \\ & & & -\mu_{n-1} & \\ & & & -\mu_{n-1} & \nu_n \end{pmatrix}$$

$$B_n = \begin{pmatrix} 0 & -\mu_1\gamma_1 & & & \\ \mu_1\gamma_1 & \ddots & & \ddots & \\ & \ddots & \ddots & & \\ & & & \ddots & -\mu_{n-1}\gamma_{n-1} \\ \mu_{n-1}\gamma_{n-1} & & & & 0 \end{pmatrix}$$

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- How about small disorder? Since the 1-D Anderson model is dynamically localized at all energies at small disorder, is \hat{M}_n as well?
- It is convenient to write \hat{M}_n in the basis $(e_1, e_{n+1}, e_2, e_{n+2}, \dots, e_n, e_{2n})$ as

$$M_n = \begin{pmatrix} \nu_1 \sigma^z & -\mu_1 S(\gamma_1) & & & \\ -\mu_1 S(\gamma_1)^t & \nu_2 \sigma^z & & \ddots & \\ & & \ddots & \ddots & \\ & & & -\mu_{n-1} S(\gamma_{n-1}) & \\ & & & -\mu_{n-1} S(\gamma_{n-1})^t & \nu_n \sigma^z \end{pmatrix}$$

where $\sigma^z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ and $S(\gamma) = \begin{pmatrix} 1 & \gamma \\ -\gamma & -1 \end{pmatrix}$

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 - $\{\nu_j\}$ are i.i.d. with nontrivial distribution ρ of compact support
 - $\gamma_j \equiv \gamma \in (0, 1) \cup (1, \infty)$ (0=isotropic, 1=Ising model)
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- With such coefficients we have

Theorem (Main Result)

For every compact interval $J \subset \mathbb{R} \setminus \{0\}$ and every $\zeta \in (0, 1)$, there exist constants $C = C(J, \zeta) < \infty$ and $\eta = \eta(J, \zeta) > 0$ such that for all $n \in \mathbb{N}$ and $j, k \in [1, n]$,

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \| P_j e^{-itM_n} \chi_J(M_n) P_k^* \| \right) \leq C e^{-\eta|j-k|\zeta}.$$

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Theorem

If there exist $\zeta \in (0, 1)$ and $C > 0, \eta > 0$ such that for all $n \in \mathbb{N}$ and $j, k \in [1, n]$,

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \|P_j e^{-itM_n} P_k^*\| \right) \leq C e^{-\eta|j-k|^\zeta}, \quad (1)$$

then for every $\varepsilon \in (0, \eta)$, there exists $C' = C'(\eta, \varepsilon, \zeta) > 0$ such that

$$\mathbb{E} \left(\sup_{t \in \mathbb{R}} \|[\tau_t^n(A), B]\| \right) \leq C' \|A\| \|B\| e^{-(\eta-\varepsilon)|j-k|^\zeta} \quad (2)$$

for all $1 \leq j < k, n \geq k, A \in \mathcal{A}_j$, and $B \in \mathcal{A}_{[k, n]}$. Furthermore, if (1) holds with $\zeta = 1$, then (2) holds with $\varepsilon = 0$.

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- Three (or more?) options to obtain an interesting result:
 - When $\gamma = 0$ (isotropic), use that (1) holds for 1-D Anderson model for arbitrary nontrivial distributions (of compact support) - the distributions do not need to be *nice*.

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 - Prove (1) from scratch, starting with regularity of Lyapunov exponents at 0.

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Proposition

Let

$$\hat{M} = \begin{pmatrix} A & B \\ -B & -A \end{pmatrix}.$$

If there exists $\lambda > 0$ such that $A \geq \lambda$ or $-A \geq \lambda$, then

$$\sigma(\hat{M}) \cap (-\lambda, \lambda) = \emptyset.$$

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- For example, if $\text{supp } \rho \subset (2, \infty)$, then the a.s. spectrum of the Anderson model is $\Sigma \subset [-2, 2] + (2, \infty)$. Thus $A_n \geq \lambda > 0$.

Proof of Main Result: A Thouless Formula

- Proof adapts strategy from [Klein/Lacroix/Speis, 1990] for Anderson model on a strip $\mathbb{Z} \times \{1, \dots, \ell\}$.

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- Lyapunov index $\gamma(E) := \frac{1}{\ell}[\gamma_1(E) + \dots + \gamma_\ell(E)]$
- Thouless formula for Anderson model on a strip:

$$\gamma(E) = \int_{\mathbb{R}} \log |E - E'| dN(E')$$

where $E \mapsto N(E)$ is the integrated density of states.

Dynamical Localization: Assumption

- Fürstenberg group $G_{\mu_E} := \overline{\langle \text{supp } \mu_E \rangle}$ is the smallest closed subgroup of the symplectic $2\ell \times 2\ell$ matrices containing $\text{supp } \mu_E$, where μ_E is the common distribution of the i.i.d. transfer matrices for the finite difference equation associated with M .

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- See [Bougerol/Lacroix, 1985] for definitions of these concepts. Suffice it to say, they generalize the notions of noncompactness and strong irreducibility required by Fürstenberg's theorem in the $\ell = 1$ case.

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- This is what we did for the anisotropic XY chain, where $\ell = 2$. We found 10 lin. indep. elements in G_{μ_E} for all $E \neq 0$. This establishes our Main Result (for the XY chain).
- Our construction follows both [Gol'dsheid/Margulis, 1989] and [Boumaza/Stolz, 2007].

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 - Representation formula for Green's function (pos. of Lyap. \Rightarrow exp. decay of Green's function)
- We establish these estimates by adapting arguments from [Klein/Lacroix/Speis, 1990] for Anderson model on a strip.
- Appropriate care must be taken to account for non-standard hopping terms.

Open Questions

- How can we remove $\chi_J(M_n)$, $J \subset \mathbb{R} \setminus \{0\}$, from our Main Result (in cases where $0 \in \Sigma_{as}$)? What happens at $E = 0$?

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- That the first Lyapunov exponent is positive is nice. The effects of the second being zero would be an interesting study.
- But to prove localization, suppose we assume $\gamma_1(0) > \gamma_2(0) > 0$ (which happens “generically”). Can one prove the required regularity of the Lyapunov exponents at 0?

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 - Main difficulty is that lack of irreducibility at 0 implies a lack of uniqueness of an invariant measure associated with the Lyapunov exponent (used to prove Hölder continuity of Lyapunov exponents).
- If we can prove it, can we extend the proof to the anisotropic XY spin chain? This would involve understanding a more abstract, higher-order dynamical system.

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Thank you!