

# Spectral Gap of $d$ -Dimensional PVBS Models

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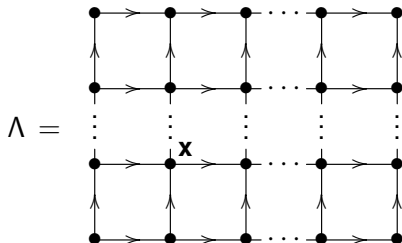
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# Hilbert Space

Let  $\Lambda$  be a finite connected subset of  $\mathbb{Z}^d$ .



The one-site Hilbert space is given by  $\mathcal{H}_x = \mathbb{C}^2$ . The Hilbert space for the whole system is given by

$$\mathcal{H}_\Lambda = \bigotimes_{x \in \Lambda} \mathcal{H}_x$$

## PVBS Hamiltonian

For each dimension  $k = 1, \dots, d$ , we assign a parameter  $\lambda_k > 0$ ,  $\lambda_k \neq 1$ .

$$H_\Lambda = \sum_{k=1}^d \sum_{\mathbf{x}, \mathbf{x}+\mathbf{e}_k \in \Lambda} h_{\mathbf{x}, \mathbf{x}+\mathbf{e}_k}$$

$$h_{\mathbf{x}, \mathbf{x}+\mathbf{e}_k} = |1, 1\rangle\langle 1, 1| + |\phi_k\rangle\langle \phi_k|, \quad \phi_k = \frac{1}{\sqrt{1 + \lambda_k^2}} (|0, 1\rangle - \lambda_k |1, 0\rangle)$$

## Ground State Space

Let  $X \subseteq \Lambda$ . Then  $\mathcal{B} = \{\psi_X^\Lambda : X \subseteq \Lambda\}$  is an orthonormal basis for  $\mathcal{H}_\Lambda$  where

$$\psi_X^\Lambda(\mathbf{x}) = \begin{cases} |1\rangle & \mathbf{x} \in X \\ |0\rangle & \mathbf{x} \notin X \end{cases}$$

Let  $\lambda^{\mathbf{x}} = \prod_{k=1}^d \lambda_k^{x_k}$ . Then ground state space for the one-species  $d$ -dimensional PVBS Hamiltonian is given by

$$\psi_\emptyset^\Lambda = \bigotimes_{\mathbf{x} \in \Lambda} |0\rangle, \quad \psi_1^\Lambda = \frac{1}{\sqrt{C_\Lambda}} \sum_{\mathbf{x} \in \Lambda} \lambda^{\mathbf{x}} \psi_{\mathbf{x}}^\Lambda \quad (1)$$

When  $\Lambda$  is the rectangular lattice with end points  $(0, 0, \dots, 0)$  and  $(n_1, n_2, \dots, n_d)$  then

$$C_\Lambda = \prod_{k=1}^d c(\lambda_k, n_k) \quad \text{where} \quad c(\lambda_k, n_k) = \sum_{i=0}^{n_k} \lambda_k^{2i}.$$

# The Spectral Gap

## Definition

Let  $\omega \in \mathcal{G}_\Gamma$  be an infinite volume ground state obtained as a weak-\* limit of finite volume ground states. Then the spectral gap of the GNS Hamiltonian  $H_\omega$  is

$$\gamma_\omega = \sup\{\delta > 0 : \text{spec}(H_\omega) \cap (0, \delta) = \emptyset\}$$

if the RHS is well defined, or zero otherwise. We say that the spectrum is gapped if  $\gamma_\omega > 0$ .

## Lower Bound for the Spectral Gap

To prove the existence of a spectral gap in the thermodynamic limit, we appeal to the following theorem.

### Theorem (Spectral Gap Estimate)

Let  $H_{\omega_0}$  be the GNS Hamiltonian of the ground state  $\omega_0 \in \mathcal{G}_{\mathbb{Z}^d}$ , and let  $\gamma_{\mathbb{Z}^d}$  be the spectral gap of  $H_{\omega_0}$ . Then

$$\gamma_{\mathbb{Z}^d} \geq \liminf_{n \geq 1} \lambda_1(n)$$

where  $\lambda_1(n)$  is the smallest nonzero eigenvalue of the frustration-free Hamiltonians  $H_{\Lambda_n}$ , where  $\Lambda_n$  is an increasing and absorbing sequence of lattices  $\Lambda_n \nearrow \mathbb{Z}^d$ .

## Conditions for the Martingale Method

The following conditions must hold for one and the same integer value  $\ell > 0$ .

- (1) There exists a constant  $d_\ell$  for which the local Hamiltonians satisfy

$$0 \leq \sum_{n=\ell}^N H_{\Lambda_n \setminus \Lambda_{n-\ell}} \leq d_\ell H_{\Lambda_N}.$$

- (2) The local Hamiltonians  $H_{\Lambda_n}$  have a non-trivial kernel  $\mathcal{G}_{\Lambda_n} \subseteq \mathcal{H}_{\Lambda_n}$ . Furthermore, there is a nonvanishing spectral gap  $\gamma_\ell > 0$  such that:

$$H_{\Lambda_n \setminus \Lambda_{n-\ell}} \geq \gamma_\ell (\mathbb{I} - G_{\Lambda_n \setminus \Lambda_{n-\ell}})$$

for all  $n \geq n_\ell$  where  $G_{\Lambda_n}$  is the orthogonal projection onto  $\mathcal{G}_{\Lambda_n}$ .

- (3) There exists a constant  $\epsilon_\ell < \frac{1}{\sqrt{\ell+1}}$  and some  $n_\ell$  such that for all  $n \geq n_\ell$ ,

$$\|G_{\Lambda_{n+1} \setminus \Lambda_{n-\ell}} E_n\| \leq \epsilon_\ell$$

where  $E_n = G_{\Lambda_n} - G_{\Lambda_{n+1}}$ .

# The Martingale Method

## Theorem

*Assume that conditions (1)-(3) are satisfied for the same integer  $\ell$ .*

*Then for any  $N$  and any  $\psi \in \mathcal{H}_{\Lambda_N}$  such that  $G_{\Lambda_N}\psi = 0$ , one has*

$$\langle \psi, H_{\Lambda_N}\psi \rangle \geq \frac{\gamma_{\ell+1}}{d_{\ell+1}} (1 - \epsilon_{\ell}\sqrt{\ell+1})^2 \|\psi\|^2 \quad (2)$$



## Sequence of Increasing Lattices

We first pick a sequence of finite volumes increasing to  $\mathbb{Z}^d$  to apply the spectral gap estimate theorem. We choose the sequence of hypercubic lattices

$$\Lambda_N = [0, N]^d \cap \mathbb{Z}^d.$$

For each term  $\Lambda_N$  of the sequence, we pick a finite sequence of lattices  $\tilde{\Lambda}_m \nearrow \Lambda_N$  to apply the martingale method. For this we choose

$$\tilde{\Lambda}_m = ([0, N]^{d-1} \times [0, m]) \cap \mathbb{Z}^d,$$

$m = 1, 2, \dots, N$ . The goal is to use the martingale method to obtain a *uniform* lower bound for the spectral gaps

$$\gamma(\Lambda_N) = \min\{\lambda \in \text{spec}(H_{\Lambda_N}) : \lambda > 0\}.$$

## Lower Bound Estimate for PVBS

### Theorem (Bounds for the Spectral Gap)

For the PVBS model defined on  $\mathbb{Z}^d$  with a single species of particle, the spectrum in the thermodynamic limit to  $\mathbb{Z}^d$  is gapped if and only if  $\lambda_k \neq 1$  for all  $k = 1, \dots, d$ . Furthermore, the spectral gap, denoted  $\gamma_{\mathbb{Z}^d}$ , is bounded by

$$\frac{\gamma(B_d)}{2^d} \prod_{k=1}^d (1 - \epsilon(\lambda_k) \sqrt{2})^2 \leq \gamma_{\mathbb{Z}^d} \leq \min \left\{ \frac{(1 - \lambda_k)^2}{1 + \lambda_k^2} : k = 1, \dots, d \right\} \quad (3)$$

where

$$\epsilon(\lambda_k) = \begin{cases} \frac{\lambda_k}{\sqrt{1 + \lambda_k^2}} & \lambda_k < 1 \\ \frac{1}{\sqrt{1 + \lambda_k^2}} & \lambda_k > 1 \end{cases} \quad (4)$$

and  $B_d$  is the  $d$ -dimensional unit hypercube.

## Sketch of Proof:

Using the rectangular sequence  $\tilde{\Lambda}_n \nearrow \Lambda_N$ , and  $\ell = 1$ , we show that

$$\|G_{\tilde{\Lambda}_{n+1} \setminus \tilde{\Lambda}_{n-1}} E_n\| = \sup_{\psi \in \mathcal{G}_{\tilde{\Lambda}_n} \cap \mathcal{G}_{\tilde{\Lambda}_{n+1}}^\perp} \frac{\|G_{\tilde{\Lambda}_{n+1} \setminus \tilde{\Lambda}_{n-1}} \psi\|}{\|\psi\|} \leq \epsilon(\lambda_d)$$

where

$$\epsilon(\lambda_d) = \begin{cases} \frac{\lambda_d}{\sqrt{1+\lambda_d^2}} & \lambda_d < 1 \\ \frac{1}{\sqrt{1+\lambda_d^2}} & \lambda_d > 1 \end{cases}$$

This satisfies  $\epsilon(\lambda_d) < \frac{1}{\sqrt{2}}$ . So by the martingale method,

$$\gamma(\Lambda_N) \geq \frac{\gamma(\Lambda_N^{(1)})}{2} (1 - \epsilon(\lambda_d)\sqrt{2})^2,$$

where

$$\Lambda_N^{(1)} = ([0, N]^{d-1} \times [0, 1]) \cap \mathbb{Z}^d.$$

## Sketch of Proof:

The sublattice  $\Lambda_N^{(1)}$  still grows as  $N \rightarrow \infty$ , so the gap could close as  $N \rightarrow \infty$ . We recursively apply the martingale method, once for each direction the lattice grows to get the estimate:

$$\gamma(\Lambda_N^{(k)}) \geq \frac{\gamma(\Lambda_N^{(k+1)})}{2} (1 - \epsilon(\lambda_{d-k})\sqrt{2})^2$$

where  $\Lambda_N^{(k)} = ([0, N]^{d-k} \times [0, 1]^k) \cap \mathbb{Z}^d$ . Since  $\Lambda_N^{(d)} = B_d$  the  $d$ -dimensional unit hypercube,

$$\gamma(\Lambda_N) \geq \frac{\gamma(B_d)}{2^d} \prod_{k=1}^d (1 - \epsilon(\lambda_k)\sqrt{2})^2$$

and the result follows from the spectral gap estimate theorem.